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Guiding properties of asymmetric hybrid plasmonic waveguides on dielectric substrates

Wei Wei, Xia Zhang*, Yongqing Huang and Xiaomin Ren

Abstract

We proposed an asymmetric hybrid plasmonic waveguide which is placed on a substrate for practical applications by introducing an asymmetry into a symmetric hybrid plasmonic waveguide. The guiding properties of the asymmetric hybrid plasmonic waveguide are investigated using finite element method. The results show that, with proper waveguide sizes, the proposed waveguide can eliminate the influence of the substrate on its guiding properties and restore its broken symmetric mode. We obtained the maximum propagation length of $2.49 \times 10^3 \ \mu m$. It is approximately equal to that of the symmetric hybrid plasmonic waveguide embedded in air cladding with comparable nanoscale confinement.

Keywords: Surface plasmons; Hybrid plasmonic waveguides; Integrated photonic devices

Background

Surface plasmons (SP) are optically induced oscillations of free electrons at the surface of a metal and can localize the guided light far beyond the diffraction limit for electromagnetic waves in dielectric media [1,2]. This could lead to miniaturized photonic circuits with a length scale much smaller than those currently achieved [3,4]. Various kinds of plasmonic waveguides including metal grooves [5,6], a chain of metal particles [7], metal stripes [8], and metal nanowires [9-11] have been proposed and investigated to realize highly integrated photonic circuits [7-12]. However, due to ohmic loss of metal [13], the propagation lengths of guided modes in plasmonic waveguides are typically short under tight confinement, which greatly limits the scope for practical applications. The main limitation of such waveguides is the trade-off between confinement and loss. Two promising approaches, the symmetric SP mode and hybrid SP mode, are proposed to optimize the balance between propagation length and mode confinement: (1) the symmetric SP mode exhibits a lower attenuation than its asymmetric counterpart, and therefore, it is sometimes referred as to long-range SP [8]; (2) in a hybrid SP mode plasmonic waveguide, the coupling between plasmonic

and waveguide modes across the gap enables 'capacitor-like' energy storage that allows subwavelength light propagation in nonmetallic regions with strong mode confinement [14]. Therefore, symmetric hybrid plasmonic (SHP) waveguides combining the two ideas of symmetric and hybrid SP modes can exhibit a quite long propagation length with strong mode confinement [15-20].

For practical implementations, an SHP waveguide needs to be placed on a substrate. The presence of the substrate breaks the symmetry of SP mode, leading to the dramatic decrease of propagation length. Here in this paper, by introducing an asymmetry into the SHP waveguide, we propose a novel asymmetric hybrid plasmonic (AHP) waveguide to eliminate the influence of a substrate on its guiding properties and restore its broken symmetric SP mode. Based on the combination of symmetric and hybrid SP modes, the AHP waveguide exhibits a quite long propagation length along with nanoscale mode confinement. In the following sections, with the finite element method (FEM), we investigate the guiding properties of the AHP waveguide on a substrate at a wavelength of 1,550 nm to target potential applications in telecommunications. Compared to an SHP waveguide with the same structure embedded in air cladding, the propagation length of the AHP waveguide is approximately the same along with a comparable normalized modal area. Moreover, the AHP waveguide has

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a horizontal slot structure featured with a horizontal low index slot, which can be convenient to be fabricated by layered deposition or thermal oxidation [21].

Methods

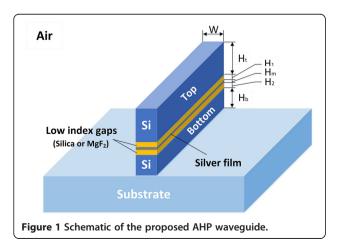
The schematic of the AHP waveguide on a silica substrate is demonstrated in Figure 1, where two layers of dielectrics (SiO₂-Si) are placed on both sides of a thin silver film. The silver film has a height of $H_{\rm m}$. The heights of the low index gaps are denoted by H_1 and H_2 , respectively. Silica is chosen to form low index gaps with respect to silicon. The heights of the top and bottom silicon layers are denoted by H_t and H_b , respectively. All these metallic and dielectric sections on the silica substrate have the same width of W. In an SHP waveguide, H_t and H_1 are equal to H_b and H_2 , respectively. However, in the AHP waveguide, $H_{\rm b}$ is smaller than $H_{\rm t}$, resulting in an asymmetry in the SHP waveguide. The optical properties of the AHP waveguide are investigated using FEM at 1,550 nm. The refractive index of silver is taken from [22]. To calculate the normalized modal area and propagation length of the AHP waveguide, we introduce Equations 1, 2, and 3 [14]:

$$A_{m} = \frac{W_{m}}{\max\{W(r)\}} = \frac{1}{\max\{W(r)\}} \iint_{\infty} W(r) d^{2}r \qquad (1)$$

where W_m is the total mode energy and W(r) is the energy density (per unit length flowed along the direction of propagation). For dispersive and lossy materials, the W(r) inside can be calculated as Equation 2:

$$W(r) = \frac{1}{2} \left(\frac{d(\varepsilon(r)\omega)}{d\omega} |E(r)|^2 + \mu_0 |H(r)|^2 \right). \tag{2}$$

The normalized modal area is defined as A_m/A_0 to quantitatively evaluate the mode confinement, where A_0



represents the diffraction-limited area in free space, $A_0 = \lambda^2/4$. The propagation length is defined as Equation 3:

$$L = \frac{1}{2\operatorname{Im}(\beta)}. (3)$$

Results and discussion

In the first section, we investigate the guiding properties and optimize structure parameters of the SHP waveguide

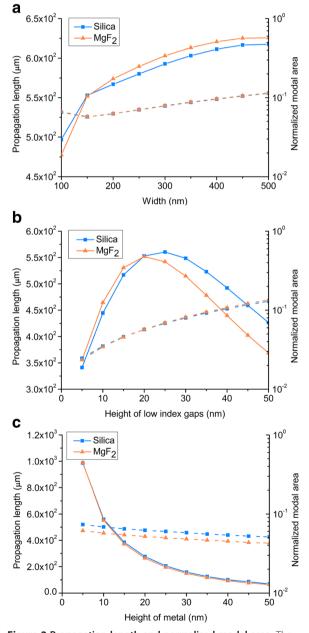


Figure 2 Propagation length and normalized modal area. They are shown versus (a) width of the waveguide, (b) height of low index gaps, and (c) height of metal stripe.

on a silica substrate via calculating the propagation length and normalized modal area. For further practical applications, the structure parameters of the SHP waveguide in the ideal condition (embedded in air cladding) are not investigated in detail here. We only compare the guiding properties between the AHP waveguide on a substrate and the SHP waveguide embedded in air cladding with the same structure parameters as the AHP waveguide. Then, in the second section, we propose the AHP waveguide by introducing an asymmetry into the SHP waveguide. Electromagnetic energy density profiles of an SHP waveguide embedded in air cladding, on a silica substrate, and an AHP waveguide on a silica substrate are demonstrated to compare SP mode distributions. We also investigate the guiding properties of the AHP waveguide as the height of mismatch varies. Here, it is worth mentioning that some values of the geometry parameters of the AHP waveguide considered in the study are reaching the limit where the local solutions of macroscopic Maxwell's equations may be not accurate enough for the descriptions of the electromagnetic properties. For more rigorous investigations, one needs to take nonlocal effects into account [14,23,24].

SHP waveguide on a substrate

Propagation length and normalized modal area are important parameters describing the mode features in a plasmonic waveguide. For applicable conditions, the SHP waveguide is always on a substrate rather than being embedded in air cladding. Therefore, in this section, we investigate the geometric dependence of the propagation length and normalized modal area of the SHP waveguide on a substrate. The propagation lengths and normalized modal areas versus the width of the low index gaps are shown in Figure 2a, where H_1 , H_2 , H_m and H_t , H_b are fixed at 20, 10, and 300 nm, respectively. The propagation lengths of silica and MgF2 increase as the width becomes wider. When the width increases, the refractive index difference brought by the substrate, which breaks the symmetric modal distribution, becomes smaller. Therefore, the propagation length increases. However, the size of waveguide increases dramatically while the propagation length increases relatively tenderly. When the width is 150 nm, there are minimum values in curves of the normalized modal area for both silica and MgF₂. At this point, the electromagnetic energy of SP mode is mostly confined in the waveguide. Due to the fact that the

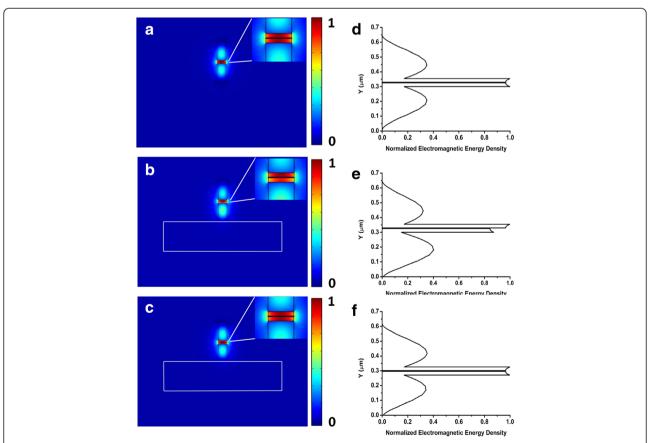


Figure 3 Electromagnetic energy density profiles of the SHP and AHP waveguides. The profiles are SHP waveguides (a) in air and (b) on a silica substrate, and (c) AHP waveguides on silica substrate. (d, e, f) Corresponding normalized electromagnetic energy densities along the *Y*-axis (from 0 to 0.6 μm) are shown.

smallest normalized modal areas are obtained at a width of 150 nm, in the following calculations, we fix the width at 150 nm. The propagation lengths and normalized modal areas versus the height of low index gaps for silica and MgF₂ are shown in Figure 2b. It is obvious that the normalized modal areas increase almost linearly with the increased heights of the low index gaps. The curves of propagation lengths are both parabolic. The propagation lengths reach the maximum values when the heights of low index gaps are equal to 25 and 20 nm, respectively. The electromagnetic energy of SP mode is mainly confined and guided in the low index gaps of the SHP waveguide. With the height of the low index gaps increasing in the rising area of the curves, more proportions of mode are confined in the gaps, which results in an extended propagation length. In this case, the mode is a hybrid mode that features both dielectric and SP characteristics [14]. With the height of the low index gaps increasing in the dropping area of the curves, the confinement becomes weaker and less proportions of mode are confined in the low index gaps, resulting in an increased loss. In the following calculations, to obtain the optimal performance of the SHP waveguide, we fix the height of low index gaps for silica and MgF₂ at 25 and 20 nm, respectively. In Figure 2c, we demonstrate the propagation lengths and normalized modal areas versus the height of metal for silica and MgF2 of the low index gaps. The propagation lengths and normalized modal areas both decrease as the height of metal increases. This can be explained as that when the height of metal becomes wider, more proportions of mode are confined in the metal, leading to increased loss and normalized modal area. Therefore, in the following, we fix the height of metal at 5 nm, emphatically considering the propagation length. Considering an ideal condition of the silica SHP waveguide being embedded in air cladding with structure parameters the same as that mentioned before, the calculated propagation length and normalized modal area are 2.38×10^3 µm and 0.076, respectively. Thus, compared to the guiding properties of the SHP embedded in air cladding, the presence of the substrate has a bad influence on the guiding properties of the SHP on the substrate. Next, we will eliminate the influence of the substrate on the guiding properties of the SHP on the substrate in an effective way.

AHP waveguide on a substrate

In this section, the structure parameters of the waveguide are the same as those in the previous section. Electromagnetic energy density profiles of the SHP waveguide in air, on a silica substrate, and an AHP waveguide on a silica substrate are shown in Figure 3a,b,c, respectively. In Figure 3a, the electromagnetic energy density profile of the SHP waveguides embedded in air cladding is symmetric. The SP mode is strongly confined and guided in

two dimensions within the low index gaps, which is bounded by the high index material and metal. However in Figure 3b, the presence of a silica substrate breaks the symmetry of the electromagnetic energy density of the SHP waveguide. The electromagnetic energy density distributes towards the upper low index gap of the SHP waveguide. When we introduce an asymmetry into the SHP waveguide on a silica substrate by decreasing $H_{\rm b}$, the asymmetric mode becomes symmetric as shown in Figure 3c. The AHP waveguide has an asymmetric structure, but its electromagnetic energy density distribution is symmetric. The asymmetric structure of the AHP waveguide restores the symmetry of the SP mode.

The height of mismatch is defined as $\Delta = H_t - H_b$ to describe the asymmetry of the AHP waveguide. The propagation length and normalized modal area of both silica and MgF₂ AHP waveguides versus the height of mismatch are shown in Figure 4, under the conditions of three different values of H_t . As shown in Figure 4a, when the height of mismatch varies from 0 to 100 nm,

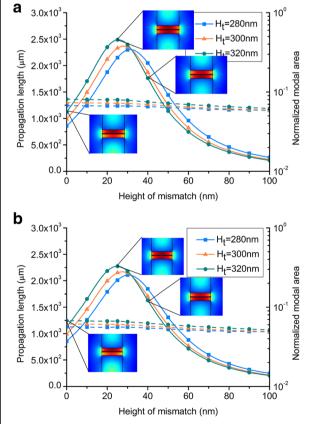


Figure 4 Propagation length and normalized modal area of silica- and MgF₂-based AHP waveguide versus height of mismatch. (a) Silica- and (b) MgF₂-based AHP waveguide. The insets indicate electromagnetic energy density profiles of different heights of mismatch.

the normalized modal area changes a little in the range of 0.06 to 0.08, which is far below the diffraction limit [25]. In a hybrid plasmonic waveguide, most proportions of the SP mode are confined in the low index gap [14]. Thus, introducing an asymmetry to the structure by varying the height of mismatch has little effect on the normalized modal area. The curves of propagation length are nearly parabolic, and the propagation length increases with the increase of H_t . As the insets of $H_t = 320$ nm as shown in Figure 4a, the electromagnetic energy of SP mode is asymmetric at $\Delta = 0$ nm. With the increase of the height of mismatch, the asymmetric mode becomes symmetric at $\Delta = 25$ nm. At this time, the propagation length reaches its maximum value. Then, with the increase of the height of mismatch, the symmetric mode becomes asymmetric again at $\Delta = 40$ nm, leading to the decrease of the propagation length. As the normalized modal areas is ultrasmall for different H_t values, we obtain the maximum propagation length of 2.49×10^3 µm for $H_t = 320$ nm. The propagation length of the AHP waveguide increases 122% than that of the SHP waveguide on a substrate. Compared to the ideal condition of the SHP in air cladding, the propagation length of the AHP waveguide is approximately equal to that of the SHP waveguide in air $(2.38 \times$ 10³ μm) with a comparable normalized modal area. Thus, the introduced asymmetry to the structure of the SHP waveguide is vital to the extension of the propagation length while exerting little effect on the normalized modal area. The phenomenon in Figure 4b is similar to that in Figure 4a, but the performance of the silica-based AHP waveguide is better than that of the MgF2-based AHP waveguide.

Conclusions

In conclusion, we reveal that the AHP waveguide combining the advantages of symmetric (long-range) SP mode and hybrid plasmonic waveguides is capable of supporting long-range propagation of the guided waves with nanoscale mode confinement. The proposed structure is realized by introducing an asymmetry into the SHP waveguide. Theoretical calculations show that the AHP waveguide can eliminate the effect of a silica substrate on the guiding properties of the SHP waveguide and restores the symmetry of SP mode. Thus, the AHP waveguide on a substrate can perform the same as the SHP waveguide embedded in air cladding. Considering different materials of the low index gaps in the AHP waveguide, the performance of the silica-based AHP waveguide is better than the MgF2-based AHP waveguide. The proposed AHP waveguide can be conveniently fabricated by existing technologies like layered deposition or thermal oxidation. This may have practical applications in highly integrated circuits as plasmonic interconnects.

Abbreviations

AHP: Asymmetric hybrid plasmonic; FEM: Finite element method; SHP: Symmetric hybrid plasmonic; SP: Surface plasmons.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

WW proposed the asymmetric idea, calculated properties of the proposed waveguide, and wrote the manuscript. XZ, YH, and XR analyzed the data and revised the manuscript. All authors read and approved the final manuscript.

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